# Majorana neutrino textures from numerical considerations. 

Bartosz Dziewit

ZTPiCE
June 13, 2011

## Introduction.

- The presently available experimental results not only suggest a non-zero neutrino mass but also constrain the patterns of neutrino masses and mixing.
- These results make it now meaningful to confront various theoretical schemes of neutrino masses with experiments.
- Neutrino mass corresponds to a Lorentz invariant renormalizable term in the Lagrangian connecting a left $\nu_{L}$ and a right-handed field $\nu_{R}$.
- Possible mass terms for neutral fermions can be written in two different ways. These are termed as Dirac and Majorana masses.


## Neutrinos in SM.

## In the Standard Model:

- there are no right-handed neutrinos $\left(\nu_{R}\right)$,
- there are only Higgs doublets of $S U(2)_{L}$,
- there are only renormalizable terms.


## Experimental arguments.

- annomalus values of solar and atmospheric neutrino fluxes,
- LSND, KamLAND, K2K.


## Theoretical predictions.

- all other fermions have mass,
- there are no symmetry principles forbiding neutrino mass terms for right handed neutrinos,


## Mass Models open problems.



## ,,Neutrino Mass Models: a road map." S.F. King arXiv:0810.0492v1 [hep-ph]

## Dirac Neutrinos.

Mass terms in SM.
All massive terms in SM occours in form:

$$
\mathcal{L}=m \bar{\psi} \Psi
$$

## Dirac Neutrinos.

Mass terms in SM.
All massive terms in SM occours in form:

$$
\mathcal{L}=m \bar{\Psi} \Psi .
$$

## Minimal extension of SM.

Intorducing right-handed neutrinos $\nu_{R}$ into the SM :

## Dirac Neutrinos.

## Mass terms in SM.

All massive terms in SM occours in form:

$$
\mathcal{L}=m \bar{\Psi} \Psi .
$$

## Minimal extension of SM.

Intorducing right-handed neutrinos $\nu_{R}$ into the SM :

$$
\mathcal{L}^{D}=-\left(\bar{\nu}_{R} m_{D} \nu_{L}+\overline{\nu_{L}} m_{D} \nu_{R}\right)+\text { h.c. }=m_{D} \bar{\nu} \nu .
$$

where: $m_{D}$ is in general $3 \times 3$ complex matrix. we can generate neutrino mass from a coupling to the Higgs

$$
\lambda_{\nu}<H>\overline{\nu_{L}} \nu_{R} \equiv m^{\nu} \bar{\nu}_{L} \nu_{R}
$$

where $<H>$ is Higgs vacuum expectation value. Physical neutrino mass of $m^{\nu} \approx 0.2[\mathrm{eV}]$ implies $\lambda_{\nu} \approx 10^{-12}$.

## Majorana Neutrinos.

## Majorana Neutrinos.

The form of a Majorana mass term is:

$$
\mathcal{L}^{M}=-\frac{1}{2} m\left(\bar{\nu}_{L} \nu_{L}^{C}+\bar{\nu}_{L}^{C} \nu_{L}\right)=-\frac{1}{2} m\left(\bar{\nu}_{L} \mathbf{C}_{L}^{T}+\text { h.c. }\right)=-\frac{1}{2} m \nu \bar{\nu} .
$$

where $\nu=\nu_{L}+\left(\nu^{C}\right)_{L}$ is a self-conjugate two-component state satisfying $\nu=\nu^{C}=C \bar{\nu}^{T}$ where $C$ is the charge conjugation matrix.

## Majorana Neutrinos.

## Majorana Neutrinos.

The form of a Majorana mass term is:

$$
\mathcal{L}^{M}=-\frac{1}{2} m\left(\bar{\nu}_{L} \nu_{L}^{C}+\bar{\nu}_{L}^{C} \nu_{L}\right)=-\frac{1}{2} m\left(\bar{\nu}_{L} \mathbf{C}_{L}^{T}+\text { h.c. }\right)=-\frac{1}{2} m \nu \bar{\nu} .
$$

where $\nu=\nu_{L}+\left(\nu^{C}\right)_{L}$ is a self-conjugate two-component state satisfying $\nu=\nu^{C}=C \bar{\nu}^{T}$ where $C$ is the charge conjugation matrix. $m$ must be generated by either an elementary Higgs triplet or by an effective operator involving two Higgs doublets arranged to transform as a triplet.

## Dirac Majorana mass terms.

## Dirac-Majorana

It is also possible to consider mixed models in which both Majorana and Dirac mass terms are present.

$$
\mathcal{L}^{D-M}=-\frac{1}{2}\left(\bar{\nu}_{L} \bar{N}_{L}^{C}\right)\left(\begin{array}{ll}
m_{T} & m_{D} \\
m_{D} & m_{S}
\end{array}\right)\binom{\nu_{R}^{C}}{N_{R}}+\text { h.c. }
$$

$m_{T}$ and $m_{S}$ are Majorana masses which transform as weak triplets and singlets, respectively while $m_{D}$ is a Dirac mass term.

## Determination of $M_{\nu}$.

## Two groups of methods:

- "top-down" method: theoretical consideration of possible textures zeros and global symmetries which seems to arise from neutrino mass matrix structure.
Analytical probes of determination $M_{\nu}$ in terms of neutrino masses values of neutrino mixing matrix elements.


## Determination of $M_{\nu}$.

## Two groups of methods:

- "top-down" method:
theoretical consideration of possible textures zeros and global symmetries which seems to arise from neutrino mass matrix structure.
Analytical probes of determination $M_{\nu}$ in terms of neutrino masses values of neutrino mixing matrix elements.
- "bottom-up" method:
relies on numerical analysis - diagonalization of many mass matrix textures leaving only these which are in agreement with present experimental data.


## Texture Zeros In The Neutrino Mass Matrix.

One-zero textures of $M_{\nu}$.

| Pattern A | Pattern B | Pattern C |
| :---: | :---: | :---: |
| $\left(\begin{array}{ccc}\mathbf{0} & \times & \times \\ \times & \times & \times \\ \times & \times & \times\end{array}\right)$ | $\left(\begin{array}{ccc}\times & \mathbf{0} & \times \\ \mathbf{0} & \times & \times \\ \times & \times & \times\end{array}\right)$ | $\left(\begin{array}{ccc}\times & \times & \mathbf{0} \\ \times & \times & \times \\ \mathbf{0} & \times & \times\end{array}\right)$ |
| Pattern D | Pattern E | Pattern F |
| $\left(\begin{array}{ccc}\times & \times & \times \\ \times & \mathbf{0} & \times \\ \times & \times & \times\end{array}\right)$ | $\left(\begin{array}{ccc}\times & \times & \times \\ \times & \times & \mathbf{0} \\ \times & \mathbf{0} & \times\end{array}\right)$ | $\left(\begin{array}{ccc}\times & \times & \times \\ \times & \times & \times \\ \times & \times & \mathbf{0}\end{array}\right)$ |

## Texture Zeros In The Neutrino Mass Matrix.

## One-zero textures of $M_{\nu}$ implications.

Note that pattern A is of particular interest, because it predicts $\langle m\rangle_{e e}=0$ (namely, the effective mass of the neutrinoless double beta decay vanishes).

$$
m_{e e}=\left(M u_{\nu}\right)_{e e}=\left(M u_{\nu}\right)_{11}=\sum U_{e i}^{2} m_{i}
$$

While:

- $\langle m\rangle_{e e} \neq 0$ must imply that neutrinos are Majorana particles,


## Texture Zeros In The Neutrino Mass Matrix.

## One-zero textures of $M_{\nu}$ implications.

Note that pattern A is of particular interest, because it predicts $\langle m\rangle_{e e}=0$ (namely, the effective mass of the neutrinoless double beta decay vanishes).

$$
m_{e e}=\left(M u_{\nu}\right)_{e e}=\left(M u_{\nu}\right)_{11}=\sum U_{e i}^{2} m_{i}
$$

While:

- $\langle m\rangle_{e e} \neq 0$ must imply that neutrinos are Majorana particles,
- $\langle m\rangle_{e e}=0$ does not necessarily imply that neutrinos are Dirac particles.


## TBM.

The lepton mixing determined from the results of neutrino experiments can be well described by the so called Tri-Bimaximal Mixing (TBM) matrix:

$$
\begin{gathered}
U_{T B M}=\left(\begin{array}{ccc}
\sqrt{\frac{2}{3}} & \frac{1}{\sqrt{3}} & 0 \\
-\frac{1}{\sqrt{6}} & \frac{1}{\sqrt{3}} & -\frac{1}{\sqrt{2}} \\
-\frac{1}{\sqrt{6}} & \frac{1}{\sqrt{3}} & \frac{1}{\sqrt{2}}
\end{array}\right) . \\
\sin ^{2} \theta_{23}=\frac{1}{2}, \quad \sin \theta_{13}=0, \quad \sin ^{2} \theta_{12}=\frac{1}{3} .
\end{gathered}
$$

In terms of the standard parameterization of lepton mixing matrix:

$$
U_{P M N S}=U_{23}\left(\theta_{23}\right) \Gamma_{\delta} U_{13}\left(\theta_{13}\right) \Gamma_{\delta}^{\star} U_{12}\left(\theta_{1} 2\right)
$$

## TBM.

For the Majorana neutrinos in the flavor basis $\left(\nu_{e}, \nu_{\mu}, \nu_{\tau}\right)$ the mass matrix which leads to the TBM mixing equals:

$$
m_{T B M}=U_{T B M} m_{\nu}^{\text {diag }} U_{T B M}^{T}
$$

In general, $m_{i}$ are complex and we can represent them as:

$$
m_{1}=\left|m_{1}\right|, \quad m_{2}=\left|m_{2}\right| e^{i 2 \phi_{2}}, \quad m_{3}=\left|m_{3}\right| e^{i 2 \phi_{3}}
$$

Here $\phi_{1}$ and $\phi_{2}$ are the Majorana CP-violating phases.

## TBM.

For the Majorana neutrinos in the flavor basis $\left(\nu_{e}, \nu_{\mu}, \nu_{\tau}\right)$ the mass matrix which leads to the TBM mixing equals:

$$
m_{T B M}=U_{T B M} m_{\nu}^{\text {diag }} U_{T B M}^{T}
$$

In general, $m_{i}$ are complex and we can represent them as:

$$
m_{1}=\left|m_{1}\right|, \quad m_{2}=\left|m_{2}\right| e^{i 2 \phi_{2}}, \quad m_{3}=\left|m_{3}\right| e^{i 2 \phi_{3}}
$$

Here $\phi_{1}$ and $\phi_{2}$ are the Majorana CP-violating phases. We can find explicitly:

$$
m_{T B M}=\left(\begin{array}{ccc}
a & b & c \\
\cdots & \frac{1}{2}(a+b+c) & \frac{1}{2}(a+b-c) \\
\cdots & \cdots & \frac{1}{2}(a+b+c)
\end{array}\right)
$$

## Large Neutrino Mixings.

## Bimaximal.

$$
U_{B M}=\left(\begin{array}{ccc}
\frac{1}{\sqrt{2}} & -\frac{1}{\sqrt{2}} & 0 \\
\frac{1}{2} & \frac{1}{2} & -\frac{1}{\sqrt{2}} \\
\frac{1}{2} & \frac{1}{2} & \frac{1}{\sqrt{2}}
\end{array}\right) \quad m_{B M}=\left(\begin{array}{ccc}
a+b & c & c \\
c & a & b \\
c & b & a
\end{array}\right)
$$

## Democratic.

$$
U_{D}=\left(\begin{array}{ccc}
\frac{1}{\sqrt{2}} & -\frac{1}{\sqrt{2}} & 0 \\
\frac{1}{\sqrt{6}} & \frac{1}{\sqrt{6}} & -\frac{2}{\sqrt{6}} \\
\frac{1}{\sqrt{3}} & \frac{1}{\sqrt{3}} & \frac{1}{\sqrt{3}}
\end{array}\right) \quad m_{D}=\left(\begin{array}{ccc}
a & b & \sqrt{2} b \\
b & a-2 c & \sqrt{2} c \\
\sqrt{2} b & \sqrt{2} c & a-c
\end{array}\right)
$$

## Texture Zeros In The Neutrino Mass Matrix.

## Two-zero textures of $M_{\nu}$.

There are totally fifteen $\left(\frac{6!}{n!(6-n)!}\right)$ possible patterns of $M_{\nu}$ with two independent vanishing entries. But only seven of them are found to be compatible with current neutrino oscillation data:

| Pattern $\mathrm{A}_{1}$ | Pattern $\mathrm{A}_{2}$ | Pattern $\mathrm{B}_{1}$ |
| :---: | :---: | :---: |
| $\left(\begin{array}{ccc}\mathbf{0} & \mathbf{0} & \times \\ \mathbf{0} & \times & \times \\ \times & \times & \times\end{array}\right)$ | $\left(\begin{array}{ccc}\mathbf{0} & \times & \mathbf{0} \\ \times & \times & \times \\ \mathbf{0} & \times & \times\end{array}\right)$ | $\left(\begin{array}{ccc}\times & \times & \mathbf{0} \\ \times & \mathbf{0} & \times \\ \mathbf{0} & \times & \times\end{array}\right)$ |
| Pattern $\mathrm{B}_{2}$ | Pattern $\mathrm{B}_{3}$ | Pattern $\mathrm{B}_{4}$ |
| $\left(\begin{array}{ccc}\times & \mathbf{0} & \times \\ \mathbf{0} & \times & \times \\ \times & \times & \mathbf{0}\end{array}\right)$ | $\left(\begin{array}{ccc}\times & \mathbf{0} & \times \\ \mathbf{0} & \mathbf{0} & \times \\ \times & \times & \times\end{array}\right)$ | $\left(\begin{array}{cc}\times & \times \\ \times & \mathbf{0} \\ \mathbf{0} & \times \\ \mathbf{0} & \times \\ \mathbf{0}\end{array}\right)$ |
| Pattern C |  |  |
| $\left(\begin{array}{ccc}\times & \times & \times \\ \times & \mathbf{0} & \times \\ \times & \times & \mathbf{0}\end{array}\right)$ |  |  |

## Texture Zeros In The Neutrino Mass Matrix.

## Three-zero textures of $M_{\nu}$ implications.

There are twenty three-zero patterns of $M_{\nu}$ which can be classified into four categories:

- Type 0 with all three diagonal matrix elements vanishing:

$$
M_{0}=\left(\begin{array}{ccc}
\mathbf{0} & \times & \times \\
\times & \mathbf{0} & \times \\
\times & \times & \mathbf{0}
\end{array}\right)
$$

- Type I with two diagonal matrix elements vanishing:

$$
M_{\mathrm{I}_{1}}=\left(\begin{array}{ccc}
\mathbf{0} & \times & \mathbf{0} \\
\times & \mathbf{0} & \times \\
\mathbf{0} & \times & \times
\end{array}\right), \quad M_{\mathrm{I}_{7}}=\left(\begin{array}{ccc}
\mathbf{0} & \mathbf{0} & \times \\
\mathbf{0} & \mathbf{0} & \times \\
\times & \times & \times
\end{array}\right)
$$

## Texture Zeros In The Neutrino Mass Matrix.

## Three-zero textures of $M_{\nu}$ implications.

There are twenty three-zero patterns of $M_{\nu}$ which can be classified into four categories:

- Type II with one diagonal matrix element vanishing:

$$
M_{\mathrm{II}_{1}}=\left(\begin{array}{ccc}
\times & \times & \mathbf{0} \\
\times & \mathbf{0} & \mathbf{0} \\
\mathbf{0} & \mathbf{0} & \times
\end{array}\right), \quad M_{\mathrm{II}_{7}}=\left(\begin{array}{ccc}
\times & \times & \mathbf{0} \\
\times & \times & \mathbf{0} \\
\mathbf{0} & \mathbf{0} & \mathbf{0}
\end{array}\right) .
$$

- Type III with three diagonal matrix elements non-vanishing:

$$
M_{\mathrm{III}}=\left(\begin{array}{ccc}
\times & \mathbf{0} & \mathbf{0} \\
\mathbf{0} & \times & \mathbf{0} \\
\mathbf{0} & \mathbf{0} & \times
\end{array}\right) .
$$

## Real Majorana Mass Matrix.

# Majorana neutrino textures from numerical considerations: the CP conserving case 

B. Dziewit, K. Kajda, J. Gluza, M. Zrałek<br>Institute of Physics, University of Silesia, Uniwersytecka 4, PL-40-007 Katowice, Poland (Dated: February 2, 2008)

Phenomenological bounds on the neutrino mixing matrix $U$ are used to determine numerically the allowed range of real elements (CP conserving case) for the symmetric neutrino mass matrix $M_{\nu}$ (Majorana case). For this purpose an adaptive Monte Carlo generator has been used. Histograms are constructed to show which forms of the neutrino mass matrix $M_{\nu}$ are possible and preferred. We confirm results found in the literature which are based on analytical calculations, though a few differences appear. These cases correspond to some textures with two zeros. The results show that actually both normal and inverted mass hierarchies are still possible at $3 \sigma$ confidence level.

## Phys. Rev. D 74 (2006) 033003 [arXiv:hep-ph/0604193].

## Real Majorana Mass Matrix.

In this paper the real symmetric $3 \times 3$ neutrino mass matrix is analyzed. It means that we assume directly the Majorana nature of neutrinos and that the investigation is restricted to the CP conserving case. A general neutrino mass matrix $M_{\nu}$ which we analyze has the following form:

$$
\left(\begin{array}{lll}
a & b & c  \tag{1}\\
b & d & e \\
c & e & f
\end{array}\right)
$$

## Real Majorana Mass Matrix.

In this paper the real symmetric $3 \times 3$ neutrino mass matrix is analyzed. It means that we assume directly the Majorana nature of neutrinos and that the investigation is restricted to the CP conserving case. A general neutrino mass matrix $M_{\nu}$ which we analyze has the following form:

$$
\left(\begin{array}{lll}
a & b & c  \tag{1}\\
b & d & e \\
c & e & f
\end{array}\right)
$$

The standard neutrino theory involves diagonalization of the neutrino mass matrix $M_{\nu}$ by use of the mixing matrix $U$ :

$$
\begin{equation*}
m_{\text {diag }}=U^{\top} M_{\nu} U \tag{2}
\end{equation*}
$$

## Real Majorana Mass Matrix.

$$
\begin{equation*}
m_{i} \leq 2.2 \mathrm{eV}, \quad\left|m_{i}-m_{j}\right|<0.05 \mathrm{eV}, \quad i, j=1,2,3 \tag{3}
\end{equation*}
$$

| $i$ | $x_{i}$ | $x_{i}^{\text {cent }}$ | $\sigma_{i}$ |
| :---: | :---: | :---: | :---: |
| 1 | $\Delta m_{32}^{2}$ | $2.6 \cdot 10^{-3}$ | $10^{-3}$ |
| 2 | $\Delta m_{21}^{2}$ | $8.3 \cdot 10^{-5}$ | $10^{-5}$ |
| 3 | $\left\|U_{e 1}\right\|$ | 0.835 | 0.045 |
| 4 | $\left\|U_{e 2}\right\|$ | 0.54 | 0.07 |
| 5 | $\left\|U_{e 3}\right\|$ | 0.1 | 0.1 |
| 6 | $\left\|U_{\mu 1}\right\|$ | 0.355 | 0.165 |
| 7 | $\left\|U_{\mu 2}\right\|$ | 0.575 | 0.155 |
| 8 | $\left\|U_{\mu 3}\right\|$ | 0.7 | 0.12 |
| 9 | $\left\|U_{\tau 1}\right\|$ | 0.365 | 0.165 |
| 10 | $\left\|U_{\tau 2}\right\|$ | 0.59 | 0.15 |
| 11 | $\left\|U_{\tau 3}\right\|$ | 0.685 | 0.125 |

The allowed absolute values of the neutrino mass squared differences $\Delta m_{32}^{2}, \Delta m_{21}^{2}$ and the allowed absolute values of the neutrino mixing matrix elements $\left|U_{i j}\right| . x_{i}^{\text {cent }}$ and $\sigma_{i}$ are the central values and the $3 \sigma$ uncertainties, respectively.

## Real Majorana Mass Matrix.

## First step - Scattering.

- Random generation of input parameters.


## Real Majorana Mass Matrix.

## First step - Scattering.

- Random generation of input parameters.
- Diagonalization of the neutrino mass matrix $M_{\nu}$.


## Real Majorana Mass Matrix.

## First step — Scattering.

- Random generation of input parameters.
- Diagonalization of the neutrino mass matrix $M_{\nu}$.
- Comparison with experimental results and saving allowed parameters

$$
\chi_{i}^{2}=\frac{\left(x_{i}^{c e n t}-x_{i}\right)^{2}}{\left(\frac{\sigma_{i}}{\alpha}\right)^{2}}
$$

## Real Majorana Mass Matrix.

## Second Step - The Adaptive Monte Carlo.

- Reading obtained points
(a) Random generation of input parameters

$$
x_{i}^{\text {cent }} \pm \xi_{i t} \delta_{i}, \quad \xi_{i t}= \begin{cases}1 & \text { it }=0 \\ 0.6 / i t & \text { it }>0\end{cases}
$$

(b) Diagonalization
(c) Comparison with experimental data and saving successive cases

## Real Majorana Mass Matrix.

## Second Step - The Adaptive Monte Carlo.

- Reading obtained points
(a) Random generation of input parameters

$$
x_{i}^{\text {cent }} \pm \xi_{i t} \delta_{i}, \quad \xi_{i t}= \begin{cases}1 & \text { it }=0 \\ 0.6 / i t & \text { it }>0\end{cases}
$$

(b) Diagonalization
(c) Comparison with experimental data and saving successive cases

- Setting new central values.

$$
\chi^{2}=\sum_{i=1}^{11} \chi_{i}^{2}
$$

## Real Majorana Mass Matrix - results.

## General case.




Figure: Three dimensional plots of allowed parameters found by the AMC procedure. On the left plot there are points obtained firstly by generating random parameters, on the right plot the points are denser as AMC looks for additional solutions in a vicinity of the parameters obtained in the first step.

## Real Majorana Mass Matrix - results.

General case - normal hierarchy.


Figure: Allowed regions of parameters for the neutrino mass matrix $M_{\nu}$ with present experimental data ( $3 \sigma$ level), the general case with normal mass hierarchy.

## Real Majorana Mass Matrix - results

## General case - normal hierarchy.








Figure: Frequency spectrum for the elements of the neutrino mass matrix $M_{\nu}$ : the general case with normal mass hierarchy.

## Real Majorana Mass Matrix - results

## General case - inverted hierarchy.




Figure: Allowed regions of parameters for the neutrino mass matrix $M_{\nu}$ with present experimental data ( $3 \sigma$ level), the general case with inverted mass hierarchy.

## Real Majorana Mass Matrix - results

## General case - inverted hierarchy.








Figure: Frequency spectrum for the elements of the neutrino mass matrix $M_{\nu}$ : the general case with inverted mass hierarchy.

## Real Majorana Mass Matrix - results

## One texture zero - A:



Figure: Allowed regions for the mass matrix with $a=0$ (A texture). The first row shows plots with $\alpha=1$ (present data, $3 \sigma$ level), the second row shows results for $\alpha=2$.

## Real Majorana Mass Matrix - results

## One texture zero - A:





Figure: Histograms of neutrino rotation angles $\sin \theta_{12}, \sin \theta_{23}$ and $\sin \theta_{13}$ for neutrino mass texture with one zero $a=0$. The histogram for $\sin \theta_{23}$ does not depend on $a$ and is the same for $a \neq 0$.

## Real Majorana Mass Matrix - results

## One texture zero - A:

| TEXTURE | ZERO | PARAMETERS | MASS RANGE | MASS MEAN | $\alpha_{0}$ |
| :---: | :---: | :---: | :---: | :---: | :---: |
| A |  | $a=0$ | $\left\|m_{3}\right\|=(0.041,0.062)$ | $<m_{3}>=0.052$ | 3.32 |
| normal |  |  | $\left\|m_{2}\right\|=(0.009,0.015)$ | $<m_{2}>=0.010$ |  |
|  |  |  | $\mid=(0.002,0.011)$ | $<m_{1}>=0.005$ |  |
|  |  |  | $\left\|m_{\beta \beta 0 \nu}\right\|=0$ | $<m_{\beta \beta 0 \nu}>=0$ |  |

Table: Masses and effective neutrinoless double beta decay mass parameter $\left\langle m_{\beta \beta 0 \nu}\right\rangle$ for allowed textures with one with one zero. The last column shows the parameter $\alpha_{0}$ for which schemes have no positive solutions.

## Real Majorana Mass Matrix - results

## Two texture zeros:

| TEXTURE | ZERO PARAMETERS | MASS RANGE | MASS MEAN | $\alpha_{0}$ |
| :---: | :---: | :---: | :---: | :---: |
| $\begin{gathered} A_{1} \text { and } A_{2} \\ \text { normal } \end{gathered}$ | $\begin{aligned} & a, b=0 \\ & a, c=0 \end{aligned}$ | $\left\|m_{3}\right\|$ $=(0.041,0.062)$ <br> $m_{2} \mid$ $=(0.009,0.015)$ <br> $\left\|m_{1}\right\|$ $=(0.002,0.012)$ <br> $\left\|m_{\beta \beta 0 \nu}\right\|=0$  | $\begin{aligned} & <m_{3}>=0.053 \\ & <m_{2}>=0.011 \\ & <m_{1}>=0.004 \\ & <m_{\beta \beta 0 \nu}>=0 \end{aligned}$ | 2.65 |
| $\begin{gathered} B_{1}, B_{2} \\ \text { degenerate } \\ \text { or } \\ \text { normal } \\ \hline \end{gathered}$ | $\begin{aligned} & c, d=0 \\ & b, f=0 \end{aligned}$ | $\left\|m_{3}\right\|$ $=(0.05,0.14)$ <br> $\left\|m_{2}\right\|$ $=(0.03,0.13)$ <br> $\left\|m_{1}\right\|=(0.02,0.13)$  <br> $\left\|m_{\beta \beta 0 \nu}\right\|=(0.02,0.13)$  | $\begin{aligned} <m_{3}> & >0.08 \\ <m_{2}> & =0.06 \\ \left\langle m_{1}>\right. & =0.06 \\ \left\langle m_{\beta \beta 0 \nu}\right. & >=0.06 \end{aligned}$ | 1.18 |
| ```B1,B2 degenerate or inverted``` | $\begin{aligned} & c, d=0 \\ & b, f=0 \end{aligned}$ | $\left\|m_{2}\right\|=(0.05,0.18)$ $\left\|m_{1}\right\|=(0.05,0.18)$ $\left\|m_{3}\right\|=(0.03,0.17)$ $\left\|m_{\beta \beta 0 \nu}\right\|=(0.05,0.18)$ | $\begin{aligned} \left\langle m_{2}\right\rangle & =0.09 \\ \left\langle m_{1}\right\rangle & =0.09 \\ \left\langle m_{3}\right\rangle & =0.07 \\ \left\langle m_{\beta \beta 0 \nu}>\right. & >=0.09 \end{aligned}$ | 1.18 |
| $\begin{gathered} \hline B_{3}, B_{4} \\ \text { degenerate } \\ \text { or } \\ \text { normal } \end{gathered}$ | $\begin{aligned} & \hline b, d=0 \\ & c, f=0 \end{aligned}$ | $\left\|m_{3}\right\|=(0.05,0.22)$ $\left\|m_{2}\right\|=(0.025,0.21)$ $\left\|m_{1}\right\|=(0.02,0.205)$ $\left\|m_{\beta \beta 0 \nu}\right\|=(0.03,0.21)$ | $\begin{aligned} \left\langle m_{3}\right\rangle & =0.08 \\ \left.<m_{2}\right\rangle & =0.06 \\ \left\langle m_{1}\right\rangle & =0.06 \\ \left\langle m_{\beta \beta 0 \nu}>\right. & >=0.06 \end{aligned}$ | 1.25 |
| $\begin{gathered} B_{3}, B_{4} \\ \text { degenerate } \\ \text { or } \\ \text { inverted } \end{gathered}$ | $\begin{aligned} & b, d=0 \\ & c, f=0 \end{aligned}$ | $\left\|m_{2}\right\|=(0.05,0.25)$ $\left\|m_{1}\right\|=(0.045,0.25)$ $\left\|m_{3}\right\|=(0.03,0.24)$ $\left\|m_{\beta \beta 0 \nu}\right\|=(0.045,0.246)$ | $\begin{aligned} <m_{2}> & =0.083 \\ <m_{1}> & =0.082 \\ <m_{3}> & =0.065 \\ <m_{\beta \beta 0 \nu}> & >0.084 \end{aligned}$ | 1.25 |
| $\begin{gathered} C \\ \text { inverted } \end{gathered}$ | $d, f=0$ | $\left\|m_{2}\right\|$ $=(0.042,0.072)$ <br> $\left\|m_{1}\right\|=(0.041,0.071)$  <br> $\left\|m_{3}\right\|$ $=(0.012,0.039)$ <br> $\left\|m_{\beta \beta 0 \nu}\right\|=(0.011,0.039)$  | $\begin{aligned} \left\langle m_{2}\right\rangle & =0.056 \\ \left.<m_{1}\right\rangle & =0.055 \\ <m_{3}> & =0.023 \\ <m_{\beta \beta 0 \nu}> & >0.022 \end{aligned}$ | 2.65 |

## Real Majorana Mass Matrix - results.

Two texture zeros.

| $i$ | $x_{i}$ | $x_{i}^{\text {cent }}$ | $\sigma_{i}$ | $A_{1}, A_{2}$ | $B_{1}-B_{4}$ | $C$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 1 | $\Delta m_{32}^{2}$ | $2.6 \cdot 10^{-3}$ | $10^{-3}$ | $2.65 \cdot 10^{-3}$ | $2.55 \cdot 10^{-3}$ | $2.61 \cdot 10^{-3}$ |
| 2 | $\Delta m_{21}^{2}$ | $8.3 \cdot 10^{-5}$ | $10^{-5}$ | $8.27 \cdot 10^{-5}$ | $8.35 \cdot 10^{-5}$ | $8.30 \cdot 10^{-5}$ |
| 3 | $\left\|U_{e 1}\right\|$ | 0.835 | 0.045 | 0.84 | 0.84 | 0.84 |
| 4 | $\left\|U_{e 2}\right\|$ | 0.54 | 0.07 | 0.54 | 0.54 | 0.54 |
| 5 | $\left\|U_{e 3}\right\|$ | 0.1 | 0.1 | 0.12 | $1.9 \cdot 10^{-3}$ | 0.06 |
| 6 | $\left\|U_{\mu 1}\right\|$ | 0.355 | 0.165 | 0.41 | 0.40 | 0.36 |
| 7 | $\left\|U_{\mu 2}\right\|$ | 0.575 | 0.155 | 0.56 | 0.58 |  |
| 8 | $\left\|U_{\mu 3}\right\|$ | 0.7 | 0.12 | 0.72 | 0.66 | 0.72 |
| 9 | $\left\|U_{\tau 1}\right\|$ | 0.365 | 0.165 | 0.36 | 0.36 | 0.41 |
| 10 | $\left\|U_{\tau 2}\right\|$ | 0.59 | 0.15 | 0.63 | 0.55 | 0.59 |
| 11 | $\left\|U_{\tau 3}\right\|$ | 0.685 | 0.125 | 0.68 | 0.75 | 0.68 |

Table: This table shows $x_{i}^{\text {cent }}$ values obtained from numerical solutions for two zero textures. It appears, that cases $A_{1}$ and $A_{2}$ coincide with solutions for one zero texture with $a=0$.

## Real Majorana Mass Matrix - results

## Two texture zeros - C:






Figure: Allowed regions for the mass matrix with $d, f=0$ (C texture). The first row shows plots with $\alpha=1$ (present data, $3 \sigma$ level). The second row shows results for $\alpha=2$.

## Real Majorana Mass Matrix - conclusions.

The most important conclusions:

- for the general case, some elements of the neutrino mass matrix $M_{\nu}$ are likely to be around zero,


## Real Majorana Mass Matrix - conclusions.

The most important conclusions:

- for the general case, some elements of the neutrino mass matrix $M_{\nu}$ are likely to be around zero,
- there are no possible numerical solutions for $M_{\nu}$ textures with number of zeros $n \geq 3$ (at $3 \sigma$ c.l.),


## Real Majorana Mass Matrix - conclusions.

The most important conclusions:

- for the general case, some elements of the neutrino mass matrix $M_{\nu}$ are likely to be around zero,
- there are no possible numerical solutions for $M_{\nu}$ textures with number of zeros $n \geq 3$ (at $3 \sigma$ c.l.),
- there are seven two zero textures which give results in agreement with present experimental data, some of them can have both normal/degenerate and inverse/degenerate mass hierarchies,


## Real Majorana Mass Matrix - conclusions.

The most important conclusions:

- for the general case, some elements of the neutrino mass matrix $M_{\nu}$ are likely to be around zero,
- there are no possible numerical solutions for $M_{\nu}$ textures with number of zeros $n \geq 3$ (at $3 \sigma$ c.l.),
- there are seven two zero textures which give results in agreement with present experimental data, some of them can have both normal/degenerate and inverse/degenerate mass hierarchies, some of them have only normal, or only inverted mass hierarchies,
- textures $B$ give small values of $\sin \theta_{13}$.


## Real Majorana Mass Matrix - conclusions.

The most important conclusions:

- for the general case, some elements of the neutrino mass matrix $M_{\nu}$ are likely to be around zero,
- there are no possible numerical solutions for $M_{\nu}$ textures with number of zeros $n \geq 3$ (at $3 \sigma$ c.l.),
- there are seven two zero textures which give results in agreement with present experimental data, some of them can have both normal/degenerate and inverse/degenerate mass hierarchies, some of them have only normal, or only inverted mass hierarchies,
- textures $B$ give small values of $\sin \theta_{13}$.
- cases with $m_{\beta \beta 0 \nu}=0$ have got only normal mass hierarchy and they all imply similar results.


## Complex Neutrino Mass Matrix.

$$
\mathcal{M}_{\nu}=W \cdot m_{\text {diag }} \cdot W^{\star}, \quad W=f \cdot U_{P M N S}^{\star} \cdot P,
$$

where:

$$
f=\left(\begin{array}{ccc}
e^{-\imath \beta_{1}} & 0 & 0 \\
0 & e^{-\imath \beta_{2}} & 0 \\
0 & 0 & e^{-\imath \beta_{3}}
\end{array}\right), P=\left(\begin{array}{ccc}
e^{-\imath \alpha_{1}} & 0 & 0 \\
0 & e^{-\imath \alpha_{2}} & 0 \\
0 & 0 & 1
\end{array}\right) .
$$

## Complex Neutrino Mass Matrix.

$$
\mathcal{M}_{\nu}=W \cdot m_{\text {diag }} \cdot W^{\star}, \quad W=f \cdot U_{P M N S}^{\star} \cdot P
$$

where:

$$
f=\left(\begin{array}{ccc}
e^{-\imath \beta_{1}} & 0 & 0 \\
0 & e^{-\imath \beta_{2}} & 0 \\
0 & 0 & e^{-\imath \beta_{3}}
\end{array}\right), P=\left(\begin{array}{ccc}
e^{-\imath \alpha_{1}} & 0 & 0 \\
0 & e^{-\imath \alpha_{2}} & 0 \\
0 & 0 & 1
\end{array}\right)
$$

We can express all terms (separatly Im and Re parts) of $\mathcal{M}_{\nu}$ as a function of:

$$
\theta_{12}, \theta_{13}, \theta_{23}, m_{1}, m_{2}, m_{3}, \delta, \alpha_{1}, \alpha_{2}, \beta_{1}, \beta_{2}, \beta_{3}
$$

and found their minimal and maximal values for current experimental data.

## Complex Neutrino Mass Matrix.






$$
\left(\begin{array}{ccc}
|a| e^{i \varphi_{a}} & |b| e^{i \varphi_{b}} & |c| e^{i \varphi_{c}} \\
|b| e^{i \varphi_{b}} & |d| e^{i{ }_{c}} & |e| e^{i i_{e}} \\
|c| e^{i \varphi_{c}} & |e| e^{i \varphi_{e}} & |b| e^{i \varphi_{f}}
\end{array}\right) \left\lvert\, \quad\left(\begin{array}{lll}
|a| e^{i \varphi_{a}} & |b| e^{i \varphi_{b}} & |c| e^{i \varphi_{c}} \\
|b| e^{i \varphi_{b}} & |d| e^{i \varphi_{d}} & |e| e^{i i_{e}} \\
|c| e^{i \varphi_{c}} & |e| e^{i \varphi_{e}} & |f| e^{i \varphi_{f}}
\end{array}\right)\right.
$$

Figure: Red scope - $\left|M_{i j}\right|, \varphi_{i j}$ with: $\delta \neq, \alpha_{i}=0, \beta_{i}=0$, blue with: $\delta \neq, \alpha_{i} \neq 0, \beta_{i}=0$.

## Complex Neutrino Mass Matrix.

## Recent work and future plans.

- non CP conserving case for real neutrino mass matrix,
- complex neutrino mass matrix:
- histograms and correspondence to real case ,
- program for automatic distinguish between possible textures,
- more general cases ... $6 \times 6$ neutrino mass matrix ...

